# On boundary value problems in three-ion electrodiffusion 

P. Amster ${ }^{\text {a,b, }, *}$, C. Rogers ${ }^{\mathrm{c}, \mathrm{d}}$<br>${ }^{\text {a }}$ Departamento de Matemática, Facultad de Ciencias Exactas y Naturales, Universidad de Buenos Aires, Buenos Aires, Argentina<br>${ }^{\text {b }}$ Consejo Nacional de Investigaciones Científicas y Técnicas (CONICET), Argentina<br>${ }^{\text {c }}$ School of Mathematics, University of New South Wales, Sydney, Australia<br>${ }^{\mathrm{d}}$ Australian Research Council Centre of Excellence for Mathematics and Statistics of Complex Systems, Australia

## Received 24 November 2006

Available online 28 March 2007
Submitted by T. Witelski


#### Abstract

The existence of solutions to a class of two-point boundary value problems in three-ion electrodiffusion is investigated via an integro-differential formulation. Boundedness by upper and lower solutions corresponding to associated boundary value problems is considered and illustrated by Painlevé II solutions of a constrained version of the original boundary value problems.


© 2007 Elsevier Inc. All rights reserved.
Keywords: Three-ion electrodiffusion; Nonlinear boundary value problems; Third order ODE's; Topological methods

## 1. Introduction

Leuchtag in [1] presented an $m$-ion electrodiffusion model consisting of the nonlinear coupled system

$$
\begin{align*}
& \frac{\mathrm{d} n_{i}}{\mathrm{~d} x}=v_{i} n_{i} p-c_{i}, \quad i=1, \ldots, m, \\
& \frac{\mathrm{~d} p}{\mathrm{~d} x}=\sum_{i=1}^{m} v_{i} n_{i}, \tag{1.1}
\end{align*}
$$

[^0]where $n_{i}$ is the number of ions with the same charge, $p$ is the electric field, the $v_{j}$ are nonzero integral signed valencies while the $c_{i}$ are real constants. In the three-ion case to be considered here, the model reduces to the following nonlinear third order equation for the electric field $p$ :
\[

$$
\begin{align*}
& p p^{\prime \prime \prime}-p^{\prime} p^{\prime \prime}-\left(v_{1}+v_{2}+v_{3}\right) p^{2} p^{\prime \prime}+\left(v_{1} v_{2}+v_{2} v_{3}+v_{3} \nu_{1}\right) p^{3} p^{\prime} \\
& \quad-\left(v_{1} c_{1}+v_{2} c_{2}+v_{3} c_{3}\right) p^{\prime}-\frac{1}{2} \nu_{1} v_{2} \nu_{3} p^{5}+v_{1} v_{2} \nu_{3}\left(c_{1}+c_{2}+c_{3}\right) x p^{3} \\
& \quad-\left[\left(v_{2}+v_{3}\right) v_{1} c_{1}+\left(v_{3}+v_{1}\right) v_{2} c_{2}+\left(v_{1}+v_{2}\right) v_{3} c_{3}\right] p^{2}=0 . \tag{1.2}
\end{align*}
$$
\]

In previous work on the model (1.1), in the two-ion case a Bäcklund transformation was applied to an underlying Painlevé II reduction [2]. Here, the three-ion model is analyzed in connection with the existence of solutions to a class of two-point boundary value problems and the construction of upper and lower solutions [3]. The work complements that initiated in [4] on the three-ion model.

## 2. A two-point boundary value problem

Let us consider the following two-point boundary value problem for (1.2):

$$
\begin{align*}
& p p^{\prime \prime \prime}-p^{\prime} p^{\prime \prime}-\left(v_{1}+v_{2}+v_{3}\right) p^{2} p^{\prime \prime}+\left(v_{1} v_{2}+v_{1} v_{3}+v_{2} \nu_{3}\right) p^{3} p^{\prime} \\
& \quad-\left(v_{1} c_{1}+v_{2} c_{2}+v_{3} c_{3}\right) p^{\prime}-\frac{1}{2} v_{1} v_{2} v_{3} p^{5}+v_{1} v_{2} v_{3}\left(c_{1}+c_{2}+c_{3}\right) x p^{3} \\
& \quad-\left[\left(v_{2}+v_{3}\right) v_{1} c_{1}+\left(v_{3}+v_{1}\right) v_{2} c_{2}+\left(v_{1}+v_{2}\right) v_{3} c_{3}\right] p^{2}=0, \\
& p(0)=p_{0}, \quad p(T)=p_{T}, \quad p^{\prime \prime}(0)=r_{0} . \tag{2.1}
\end{align*}
$$

If we set $u=p^{\prime \prime} / p, u_{0}=r_{0} / p_{0}$, then for $p \neq 0$ this boundary value problem is equivalent to:

$$
\begin{align*}
p^{2} u^{\prime} & -\left(v_{1}+v_{2}+v_{3}\right) p^{2} p^{\prime \prime}+\left(v_{1} v_{2}+v_{1} v_{3}+v_{2} v_{3}\right) p^{3} p^{\prime} \\
& \quad-\left(v_{1} c_{1}+v_{2} c_{2}+v_{3} c_{3}\right) p^{\prime}-\frac{1}{2} v_{1} v_{2} v_{3} p^{5}+v_{1} v_{2} v_{3}\left(c_{1}+c_{2}+c_{3}\right) x p^{3} \\
= & {\left[\left(v_{2}+v_{3}\right) v_{1} c_{1}+\left(v_{3}+v_{1}\right) v_{2} c_{2}+\left(v_{1}+v_{2}\right) v_{3} c_{3}\right] p^{2}, } \\
p^{\prime \prime}= & p u, \quad p(0)=p_{0}, \quad p(T)=p_{T}, \quad u(0)=u_{0} . \tag{2.2}
\end{align*}
$$

This in turn, is equivalent to the second order integro-differential boundary value problem:

$$
\begin{align*}
& p^{\prime \prime}=-C_{3}+\varphi(x) p+C_{1}\left(p^{\prime}-p^{\prime}(0)\right) p+\frac{C_{2}}{2} p^{3} p \int_{0}^{x}\left(C_{4} p^{2}(t)+C_{5} t\right) p(t) \mathrm{d} t \\
& p(0)=p_{0}, \quad p(T)=p_{T}, \tag{2.3}
\end{align*}
$$

where

$$
\begin{aligned}
& C_{1}=v_{1}+v_{2}+v_{3}, \quad C_{2}=-\left(v_{1} v_{2}+v_{1} v_{3}+v_{2} v_{3}\right), \quad C_{3}=v_{1} c_{1}+v_{2} c_{2}+v_{3} c_{3}, \\
& C_{4}=\frac{1}{2} v_{1} v_{2} v_{3}, \quad C_{5}=-v_{1} v_{2} v_{3}\left(c_{1}+c_{2}+c_{3}\right), \\
& C_{6}=\left(v_{2}+v_{3}\right) v_{1} c_{1}+\left(v_{3}+v_{1}\right) \nu_{2} c_{2}+\left(v_{1}+v_{2}\right) v_{3} c_{3},
\end{aligned}
$$

and

$$
\varphi(x)=u_{0}-\frac{1}{2} C_{2} p_{0}^{2}+\frac{C_{3}}{p_{0}}+C_{6} x .
$$

Here, we shall assume that the associated linear operator $\mathcal{L} p:=p^{\prime \prime}-\varphi p$ is below resonance for the Dirichlet conditions, that is

$$
\begin{equation*}
\varphi_{i}:=\inf _{0 \leqslant x \leqslant T} \varphi(x)=\frac{r_{0}+C_{3}}{p_{0}}-C_{2} \frac{p_{0}^{2}}{2}+\min \left\{0, C_{6} T\right\}>-\left(\frac{\pi}{T}\right)^{2} \tag{2.4}
\end{equation*}
$$

By standard results, condition (2.4) implies that $\mathcal{L}: H^{2} \cap H_{0}^{1}(0, T) \rightarrow L^{2}(0, T)$ is invertible. In particular, a straightforward computation shows that if $p \in H^{2} \cap H_{0}^{1}(0, T)$ then

$$
\|p\|_{L^{2}} \leqslant \frac{1}{\left(\frac{\pi}{T}\right)^{2}+\varphi_{i}}\|\mathcal{L} p\|_{L^{2}}
$$

and

$$
\left\|p^{\prime}\right\|_{L^{2}} \leqslant \frac{\frac{\pi}{T}}{\left(\frac{\pi}{T}\right)^{2}+\varphi_{i}}\|\mathcal{L} p\|_{L^{2}}
$$

Furthermore, writing $p^{\prime \prime}=\mathcal{L} p+\varphi p$, we also deduce that

$$
\left\|p^{\prime \prime}\right\|_{L^{2}} \leqslant(1+k)\|\mathcal{L} p\|_{L^{2}}
$$

where

$$
k=\frac{\max \left\{\left|\frac{r_{0}+C_{3}}{p_{0}}-C_{2} \frac{p_{0}^{2}}{2}\right|,\left|\frac{r_{0}+C_{3}}{p_{0}}-C_{2} \frac{p_{0}^{2}}{2}+C_{6} T\right|\right\}}{\left(\frac{\pi}{T}\right)^{2}+\varphi_{i}} .
$$

Hence,

$$
\left\|p^{\prime}\right\|_{C} \leqslant T^{1 / 2}\left\|p^{\prime \prime}\right\|_{L^{2}} \leqslant T^{1 / 2}(1+k)\|\mathcal{L} p\|_{L^{2}}
$$

and setting

$$
N=T^{1 / 2} \max \left\{\frac{\frac{\pi}{T}}{\left(\frac{\pi}{T}\right)^{2}+\varphi_{i}}, 1+k\right\}
$$

we conclude that

$$
\|p\|_{C^{1}}:=\max \left\{\|p\|_{C},\left\|p^{\prime}\right\|_{C}\right\} \leqslant N\|\mathcal{L} p\|_{L^{2}} .
$$

In order to establish a sufficient condition for the existence of solutions, let us define a polynomial $\gamma: \mathbb{R} \rightarrow \mathbb{R}$ given by

$$
\begin{equation*}
\gamma(M)=2\left|C_{1}\right| M^{2}+\frac{\left|C_{2}\right|}{2} M^{3}+T M^{2}\left(\left|C_{4}\right| M^{2}+\left|C_{5}\right| T\right) . \tag{2.5}
\end{equation*}
$$

From the limits

$$
\lim _{M \rightarrow 0^{+}} \frac{\gamma(M)}{M}=0, \quad \lim _{M \rightarrow+\infty} \frac{\gamma(M)}{M}=+\infty
$$

it is clear that the function $M-N T^{1 / 2} \gamma(M)$ achieves a positive maximum $A_{\max }$ at some value $M_{\text {max }}>0$.

The following result may be established:
Theorem 2.1. Assume that (2.4) holds, and that $\Theta \leqslant A_{\max }$, where

$$
\begin{equation*}
\Theta:=\max \left\{\left|p_{0}\right|,\left|p_{T}\right|, \frac{\left|p_{T}-p_{0}\right|}{T}\right\}+N\left(\left|C_{3}\right| T^{1 / 2}+\|\varphi B\|_{L^{2}}\right) \tag{2.6}
\end{equation*}
$$

and

$$
B(t)=\left(\frac{p_{T}-p_{0}}{T}\right) t+p_{0} .
$$

Then the boundary value problem (2.1) admits at least one classical solution.

Proof. Let us define an operator $\mathcal{T}: C^{1}[0, T] \rightarrow C^{1}[0, T]$ given by $\mathcal{T} q=p$, where $p$ is the unique solution of the linear problem

$$
\left\{\begin{array}{l}
p^{\prime \prime}(x)-\varphi(x) p(x)=V(q)-C_{3}, \\
p(0)=p_{0}, \quad p(T)=p_{T},
\end{array}\right.
$$

with

$$
V(q)(x):=C_{1}\left(q^{\prime}(x)-q^{\prime}(0)\right) q(x)+\frac{C_{2}}{2} q^{3}(x)+q(x) \int_{0}^{x}\left(C_{4} q^{2}(t)+C_{5} t\right) q(t) \mathrm{d} t .
$$

By standard results, $\mathcal{T}$ is well defined and compact. Moreover, if $\theta$ is the unique function satisfying

$$
\mathcal{L} \theta=-C_{3}, \quad \theta(0)=p_{0}, \quad \theta(T)=p_{T},
$$

then

$$
\|\theta-B\|_{C^{1}} \leqslant N\left\|C_{3}-\varphi B\right\|_{L^{2}} \leqslant N\left(\left|C_{3}\right| T^{1 / 2}+\|\varphi B\|_{L^{2}}\right)
$$

Hence, $\|\theta\|_{C^{1}} \leqslant \Theta$, where $\Theta$ is given by (2.6). Moreover, the following bound is obtained for $p=\mathcal{T} q$ :

$$
\|p-\theta\|_{C^{1}} \leqslant N\|\mathcal{L}(p-\theta)\|_{L^{2}}=N\|V(q)\|_{L^{2}} .
$$

It is readily shown that if $\|q\|_{C^{1}} \leqslant M$ then $|V(q)(x)| \leqslant \gamma(M)$, with $\gamma(M)$ as in (2.5). Thus, $\|p\|_{C^{1}} \leqslant \Theta+N T^{1 / 2} \gamma(M)$, and since $\Theta \leqslant A_{\max }$ we conclude that if $\|q\|_{C^{1}} \leqslant M_{\max }$ then $\|p\|_{C^{1}} \leqslant M_{\max }$. The result now follows from Schauder's Fixed Point Theorem.

As a particular consequence we obtain:

Corollary 2.1. Assume that $\min \left\{0, C_{6} T\right\}>L-\left(\frac{\pi}{T}\right)^{2}$ for some $L>0$. Then there exist constants $\delta_{0}, \delta_{T}>0$ such that the boundary value problem (2.1) admits at least one classical solution for any choice of the parameters $p_{0}, p_{T}, C_{3}$ and $r_{0}$ satisfying:

$$
\left|p_{T}\right|<\delta_{T}, \quad \frac{\left|r_{0}\right|+\left|C_{3}\right|}{L}<\left|p_{0}\right|<\delta_{0} .
$$

Proof. From the hypotheses, it is clear that if $C_{2}$ and $C_{6}$ are fixed then (2.4) holds for $\delta_{0}$ small enough. Furthermore, it is observed that $\|\varphi\|_{C}$ and the constant $N$ remain bounded as $p_{0} \rightarrow 0$, provided $\left|\frac{r_{0}+C_{3}}{p_{0}}\right|<L$. On the other hand, if $p_{0}, p_{T} \rightarrow 0$ then $\|B\|_{C^{1}} \rightarrow 0$ and $C_{3} \rightarrow 0$. Hence $\Theta \rightarrow 0$, and the result follows.

## 3. The case $v_{1}+v_{2}+v_{3}=0$. Upper and lower solutions

Here we focus on the boundary value problem

$$
\begin{align*}
& p^{\prime \prime}=-C_{3}+\varphi(x) p+\frac{C_{2}}{2} p^{3}+p \int_{0}^{x}\left(C_{4} p^{2}(t)+C_{5} t\right) p(t) \mathrm{d} t \\
& p(0)=p_{0}, \quad p(T)=p_{T} \tag{3.1}
\end{align*}
$$

corresponding to the case $C_{1}=\sum v_{i}=0$. Note that this condition implies that $C_{2} \geqslant 0$.
The quantities $\alpha$ and $\beta$ are termed lower and upper solutions respectively for the boundary value problem (3.1) if

$$
\begin{align*}
& \alpha^{\prime \prime} \geqslant-C_{3}+\varphi(x) \alpha+\frac{C_{2}}{2} \alpha^{3}+\alpha \int_{0}^{x}\left(C_{4} \alpha^{2}(t)+C_{5} t\right) \alpha(t) \mathrm{d} t  \tag{3.2}\\
& \beta^{\prime \prime} \leqslant-C_{3}+\varphi(x) \beta+\frac{C_{2}}{2} \beta^{3}+\beta \int_{0}^{x}\left(C_{4} \beta^{2}(t)+C_{5} t\right) \beta(t) \mathrm{d} t, \tag{3.3}
\end{align*}
$$

and

$$
\begin{equation*}
\alpha(0) \leqslant p_{0} \leqslant \beta(0), \quad \alpha(T) \leqslant p_{T} \leqslant \beta(T) . \tag{3.4}
\end{equation*}
$$

Then we have:
Theorem 3.1. Let $\alpha$ and $\beta$ be respectively lower and upper solutions of the boundary value problem (3.1), with $0 \leqslant \alpha \leqslant \beta$, and assume that $\nu_{1} \nu_{2}\left(\nu_{1}+\nu_{2}\right)>0, c_{1}+c_{2}+c_{3} \leqslant 0$. Then (3.1) admits a solution $u$ with $\alpha \leqslant u \leqslant \beta$.

Proof. From the assumptions, it readily follows that if $0 \leqslant p \leqslant q$ then

$$
p(x) \int_{0}^{x}\left(C_{4} p^{2}(t)+C_{5} t\right) p(t) \mathrm{d} t \geqslant q(x) \int_{0}^{x}\left(C_{4} q^{2}(t)+C_{5} t\right) q(t) \mathrm{d} t .
$$

Let us fix a nonnegative constant $\lambda$ such that

$$
\lambda \geqslant \frac{3}{2} C_{2}\|\beta\|_{C}^{2}+\varphi_{s}
$$

with

$$
\varphi_{s}:=\sup _{0 \leqslant x \leqslant T} \varphi(x)=\frac{r_{0}+C_{3}}{p_{0}}-C_{2} \frac{p_{0}^{2}}{2}+\max \left\{0, C_{6} T\right\} .
$$

Then the function $\varphi p+\frac{C_{2}}{2} p^{3}-\lambda p$ is nonincreasing in $p$ for $\alpha(x) \leqslant p \leqslant \beta(x)$. Next, define a fixed point operator $\mathcal{T}$ such that for fixed $q \in C[0, T], p=\mathcal{T} q$ is the unique solution of the linear problem

$$
\begin{aligned}
& p^{\prime \prime}-\lambda p=-C_{3}+\varphi q+\frac{C_{2}}{2} q^{3}-\lambda q+q \int_{0}^{x}\left(C_{4} q^{2}(t)+C_{5} t\right) q(t) \mathrm{d} t, \\
& p(0)=p_{0}, \quad p(T)=p_{T} .
\end{aligned}
$$

By standard results, $\mathcal{T}: C[0, T] \rightarrow C[0, T]$ is well defined and compact. Moreover, if $\alpha \leqslant q \leqslant \beta$, for $p=\mathcal{T} q$ it follows that

$$
p^{\prime \prime}-\lambda p \leqslant-C_{3}+\varphi \alpha+\frac{C_{2}}{2} \alpha^{3}-\lambda \alpha+\alpha \int_{0}^{x}\left(C_{4} \alpha^{2}(t)+C_{5} t\right) \alpha(t) \mathrm{d} t \leqslant \alpha^{\prime \prime}-\lambda \alpha
$$

and

$$
p^{\prime \prime}-\lambda p \geqslant-C_{3}+\varphi \beta+\frac{C_{2}}{2} \beta^{3}-\lambda \beta+\beta \int_{0}^{x}\left(C_{4} \beta^{2}(t)+C_{5} t\right) \beta(t) \mathrm{d} t \geqslant \beta^{\prime \prime}-\lambda \beta
$$

From (3.4) and the maximum principle we conclude that $\alpha \leqslant p \leqslant \beta$.
The result follows by applying Schauder's Theorem to the bounded, convex and closed set $\{u \in C[0, T]: \alpha \leqslant u \leqslant \beta\}$.

An analogous result when $\alpha \leqslant \beta \leqslant 0$ may be established, namely:
Theorem 3.2. Let $\alpha$ and $\beta$ be respectively lower and an upper solutions of (3.1), with $\alpha \leqslant \beta \leqslant 0$, and assume that $\nu_{1} v_{2}\left(v_{1}+v_{2}\right)<0, c_{1}+c_{2}+c_{3} \leqslant 0$. Then the boundary value problem (3.1) admits a solution $u$ with $\alpha \leqslant u \leqslant \beta$.

## 4. An iterative quasilinearization method

Here, the existence of an ordered couple of lower and upper solutions is assumed and an iterative scheme that converges to a solution of problem (3.1) is constructed. Under appropriate conditions, the convergence is proved to be quadratic.

In this connection, it proves convenient to write (3.1) $)_{1}$ as

$$
p^{\prime \prime}-\frac{C_{2}}{2} p^{3}=F(p)
$$

where the mapping $F: C[0,1] \rightarrow C[0,1]$ is given by

$$
\begin{equation*}
F(p)=-C_{3}+\varphi p+p \int_{0}^{x}\left(C_{4} p^{2}(t)+C_{5} t\right) p(t) \mathrm{d} t \tag{4.1}
\end{equation*}
$$

It is readily seen that $F$ is infinitely Fréchet differentiable, with

$$
\begin{equation*}
D F(p)[q]=\varphi q+p \int_{0}^{x}\left(3 C_{4} p^{2}(t)+C_{5} t\right) q(t) \mathrm{d} t+q \int_{0}^{x}\left(C_{4} p^{2}(t)+C_{5} t\right) p(t) \mathrm{d} t \tag{4.2}
\end{equation*}
$$

and

$$
\begin{align*}
D^{2} F(p)[q, r]= & 6 C_{4} p \int_{0}^{x} p q r \mathrm{~d} t+q \int_{0}^{x}\left(3 C_{4} p^{2}(t)+C_{5} t\right) r(t) \mathrm{d} t \\
& +r \int_{0}^{x}\left(3 C_{4} p^{2}(t)+C_{5} t\right) q(t) \mathrm{d} t \tag{4.3}
\end{align*}
$$

As a preliminary we note appropriate comparison and existence-uniqueness results for the semilinear operator $S_{\lambda} p:=p^{\prime \prime}-\frac{C_{2}}{2} p^{3}-\lambda p$ with $\lambda, C_{2} \geqslant 0$ (proofs are straightforward):

Lemma 4.1. If $p, q \in H^{2}(0, T)$ satisfy

$$
S_{\lambda} p \geqslant S_{\lambda} q \quad \text { a.e. } \quad p(0) \leqslant q(0), \quad p(T) \leqslant q(T)
$$

then $p \leqslant q$.
Lemma 4.2. Let $\xi \in L^{2}(0, T)$ and $p_{0}, p_{T} \in \mathbb{R}$. Then the boundary value problem

$$
S_{\lambda} p=\xi, \quad p(0)=p_{0}, \quad p(T)=p_{T}
$$

admits a unique solution $p \in H^{2}(0, T)$.

In what follows, we shall consider only the situation when Theorem 3.1 applies. Analogous conclusions hold when Theorem 3.2 obtains.

In order to construct an iterative Newton-type scheme for the boundary value problem (3.1), the following result is required.

Lemma 4.3. Let the assumptions of Theorem 3.1 hold. Then there exists $p \in C^{2}[0, T]$ such that $\alpha \leqslant p \leqslant \beta$, and

$$
\left\{\begin{array}{l}
S_{0} p=F(\alpha)+D F(\alpha)(p-\alpha)  \tag{4.4}\\
p(0)=p_{0}, \quad p(T)=p_{T}
\end{array}\right.
$$

Proof. Fix a nonnegative constant $\lambda \geqslant \varphi_{s}$, with $\varphi_{s}$ as in Theorem 3.1, and define a function $\Pi:[0, T] \times \mathbb{R} \rightarrow \mathbb{R}$ by

$$
\Pi(x, p)= \begin{cases}p & \text { if } \alpha(x) \leqslant p \leqslant \beta(x) \\ \alpha(x) & \text { if } p<\alpha(x) \\ \beta(x) & \text { if } p>\beta(x)\end{cases}
$$

From Lemma 4.2 and Schauder's Theorem, it is seen that the following quasilinear problem admits at least one solution:

$$
\left\{\begin{array}{l}
S_{\lambda} p=F(\alpha)+D F(\alpha)(\Pi(p)-\alpha)-\lambda \Pi(p) \\
p(0)=p_{0}, \quad p(T)=p_{T}
\end{array}\right.
$$

We now write $S_{\lambda} p=F(\alpha)-\lambda \alpha+(D F(\alpha)-\lambda I)(\Pi(p)-\alpha)$, and from the choice of $\lambda$ it follows immediately that $S_{\lambda} p \leqslant F(\alpha)-\lambda \alpha \leqslant S_{\lambda} \alpha$.

On the other hand, consider the Taylor expansion

$$
F(z)=F(\alpha)+D F(\alpha)(z-\alpha)+R(z)
$$

From the fact that $D^{2} F(p)[q, r] \leqslant 0$ for $p, q, r \geqslant 0$, it follows that $R(z) \leqslant 0$ for $z \geqslant \alpha$. Then

$$
S_{\lambda} p=F(\Pi(p))-R(\Pi(p))-\lambda \Pi(p) \geqslant F(\beta)-\lambda \beta \geqslant S_{\lambda} \beta .
$$

Hence, $\alpha \leqslant p \leqslant \beta$, and $S_{0} p=F(\alpha)+D F(\alpha)(p-\alpha)$. Again, we may write $S_{0} p=F(p)-R(p)$, where the Taylor remainder $R(p)$ is nonpositive, and conclude that $S_{0} p \geqslant F(p)$. Thus $p$ is a lower solution of the boundary value problem (3.1).

Next, we define a sequence as follows. Start with $p_{1}=\alpha$, then from Lemma 4.2 we may choose a lower solution $p_{2}$ with $p_{1} \leqslant p_{2} \leqslant \beta$ satisfying (4.4). Iteration of this process produces a nondecreasing sequence

$$
p_{1} \leqslant p_{2} \leqslant p_{3} \leqslant \cdots \leqslant \beta
$$

such that

$$
\begin{equation*}
S_{0} p_{n+1}=F\left(p_{n}\right)+D F\left(p_{n}\right)\left(p_{n+1}-p_{n}\right), \quad p_{n+1}(0)=p_{0}, \quad p_{n+1}(T)=p_{T}, \tag{4.5}
\end{equation*}
$$

where $\left\{p_{n}\right\}$ converges pointwise to some function $p$. From standard results (Dini's Theorem), $p_{n} \rightarrow p$ uniformly, and use of (4.5) shows that $p$ is a solution of (3.1). In order to prove the quadratic convergence of $\left\{p_{n}\right\}$ we impose an extra requirement:

Theorem 4.4. Let the assumptions of Theorem 3.1 hold. Further, assume that

$$
\begin{equation*}
\mu+\varphi+\frac{C_{2}}{2}\left(p_{n+1}^{2}+p_{n+1} p_{n}+p_{n}^{2}\right)+\int_{0}^{x}\left(C_{4} \beta^{2}+C_{5} t\right) \beta \mathrm{d} t \geqslant 0 \tag{4.6}
\end{equation*}
$$

for some $n \geqslant 2$ and some constant $\mu<\left(\frac{\pi}{T}\right)^{2}$, and that $k_{0} k_{1}<\frac{1}{T}$, where

$$
k_{0}=\frac{\frac{\pi}{T}}{\left(\frac{\pi}{T}\right)^{2}-\mu}, \quad k_{1}=\left\|\beta \int_{0}^{x}\left(3 C_{4} \beta^{2}+C_{5} t\right) \mathrm{d} t\right\|_{C}
$$

Then

$$
\left\|p_{n+1}-p_{n}\right\|_{C} \leqslant k\left\|p_{n}-p_{n-1}\right\|_{C}^{2}
$$

for some constant $k$ independent of $n$. In particular, the sequence defined by (4.6) converges quadratically to a solution of the boundary value problem (3.1).

Proof. In the context the previous proof, define $E_{n}=p_{n}-p_{n-1}$. Then $\left\{E_{n}\right\}$ is pointwise nonincreasing and tends to 0 as $n \rightarrow \infty$. Moreover, for $n \geqslant 2$

$$
S_{0} p_{n+1}-S_{0} p_{n}=R_{n}+D F\left(p_{n}\right)\left(p_{n+1}-p_{n}\right)
$$

where $R_{n}:=F\left(p_{n}\right)-F\left(p_{n-1}\right)-D F\left(p_{n-1}\right)\left(p_{n}-p_{n-1}\right)$ is the Taylor remainder. It follows that

$$
\begin{aligned}
E_{n+1}^{\prime \prime}+\mu E_{n+1}= & R_{n}+p_{n} \int_{0}^{x}\left(3 C_{4} p_{n}^{2}+C_{5} t\right) E_{n+1} \mathrm{~d} t \\
& +\left(\mu+\varphi+\frac{C_{2}}{2}\left(p_{n+1}^{2}+p_{n+1} p_{n}+p_{n}^{2}\right)+\int_{0}^{x}\left(C_{4} p_{n}^{2}+C_{5} t\right) p_{n} \mathrm{~d} t\right) E_{n+1} \\
\geqslant & R_{n}+p_{n} \int_{0}^{x}\left(3 C_{4} p_{n}^{2}+C_{5} t\right) E_{n+1} \mathrm{~d} t
\end{aligned}
$$

Thus, if $\Phi \in H^{2} \cap H_{0}^{1}(0, T)$ denotes the unique solution of

$$
\begin{equation*}
\Phi^{\prime \prime}+\mu \Phi=R_{n}+p_{n} \int_{0}^{x}\left(3 C_{4} p_{n}^{2}+C_{5} t\right) E_{n+1} \mathrm{~d} t \tag{4.7}
\end{equation*}
$$

then the standard comparison principle implies that $E_{n+1} \leqslant \Phi$. Moreover,

$$
\|\Phi\|_{C} \leqslant T k_{0}\left\|\Phi^{\prime \prime}+\mu \Phi\right\|_{C} \leqslant T k_{0}\left(\left\|R_{n}\right\|_{C}+k_{1}\left\|E_{n+1}\right\|_{C}\right)
$$

and hence

$$
\left\|E_{n+1}\right\|_{C} \leqslant \frac{T k_{0}}{1-T k_{0} k_{1}}\left\|R_{n}\right\|_{C}
$$

Furthermore, from the Taylor expansion of $F$ we deduce that

$$
\left\|R_{n}\right\|_{C} \leqslant T\left(6\left|C_{4}\right|\|\beta\|_{C}^{2}+\left|C_{5}\right| \frac{T}{2}\right)\left\|E_{n}\right\|_{C}^{2}
$$

Finally, note that if (4.6) holds, then it also holds for any $m \geqslant n$, and the result is established.

## 5. Painlevé II

Here, we present Painlevé II solutions of a suitably constrained version of boundary value problem (2.1). Thus, we obtain particular nonconstant solutions of (2.1) such that

$$
\begin{equation*}
p^{\prime \prime}+\lambda p^{3}+\mu x p+v=0 \tag{5.1}
\end{equation*}
$$

for appropriate constants $\lambda, \mu$ and $\nu$. Note that if $\lambda$ is negative, the transformation $p(x) \mapsto$ $Y(x):=\sigma p(\omega x)$ for suitable choice of $\sigma$ and $\omega$ gives a solution of the standard Painlevé II equation $Y^{\prime \prime}=2 Y^{3} \pm x Y+C$. One use of (5.1) to eliminate $p^{\prime \prime \prime}$ and $p^{\prime \prime}$ in (2.1), it is seen that

$$
\begin{aligned}
& -p\left[3 \lambda p^{2} p^{\prime}+\mu p+\mu x p^{\prime}\right]+p^{\prime}\left[\lambda p^{3}+\mu x p+\nu\right]+\left(v_{1}+\nu_{2}+\nu_{3}\right) p^{2}\left[\lambda p^{3}+\mu x p+\nu\right] \\
& +\left(\nu_{1} \nu_{2}+\nu_{1} \nu_{3}+\nu_{2} \nu_{3}\right) p^{3} p^{\prime}-\left(\nu_{1} c_{1}+\nu_{2} c_{2}+\nu_{3} c_{3}\right) p^{\prime}-\frac{1}{2} \nu_{1} \nu_{2} \nu_{3} p^{5} \\
& +\nu_{1} \nu_{2} \nu_{3}\left(c_{1}+c_{2}+c_{3}\right) x p^{3}-\left[\left(\nu_{2}+\nu_{3}\right) \nu_{1} c_{1}+\left(\nu_{1}+\nu_{3}\right) \nu_{2} c_{2}+\left(\nu_{1}+\nu_{2}\right) \nu_{3} c_{3}\right] p^{2} \\
& =0,
\end{aligned}
$$

whence we obtain:

$$
\begin{align*}
& \lambda= \frac{1}{2}\left(\nu_{2} v_{3}+v_{1} v_{3}+v_{1} v_{2}\right) \\
& \mu=\left(v_{1} c_{1}+v_{2} c_{2}+v_{3} c_{3}\right)\left(\nu_{1}+v_{2}+v_{3}\right) \\
&-\left[\left(v_{2}+v_{3}\right) v_{1} c_{1}+\left(v_{1}+v_{3}\right) v_{2} c_{2}+\left(v_{1}+v_{2}\right) v_{3} c_{3}\right] \\
& \nu= v_{1} c_{1}+v_{2} c_{2}+v_{3} c_{3} \\
& \mu\left(v_{1}+v_{2}+v_{3}\right)+v_{1} v_{2} v_{3}\left(c_{1}+c_{2}+c_{3}\right)=0, \\
&\left(v_{1}+v_{2}\right)\left(v_{1}+v_{3}\right)\left(v_{2}+v_{3}\right)=0 \tag{5.2}
\end{align*}
$$

In view of the latter condition, we proceed with the constraint $\nu_{1}+\nu_{2}=0, \nu_{3} \neq 0$ whence $c_{3}=0$, and

$$
\begin{equation*}
\lambda=-v_{1}^{2} / 2, \quad \mu=v_{1}^{2}\left(c_{1}+c_{2}\right), \quad \nu=c_{1} v_{1}+c_{2} v_{2}=v_{1}\left(c_{1}-c_{2}\right) . \tag{5.3}
\end{equation*}
$$

Finally, from the boundary conditions in (2.1) we obtain the constraint

$$
\begin{equation*}
r_{0}=\frac{v_{1}^{2}}{2} p_{0}^{3}-v_{1}\left(c_{1}-c_{2}\right) . \tag{5.4}
\end{equation*}
$$

Analogous results hold by cyclic interchange for $\left\{\nu_{2}+\nu_{3}=0, \nu_{1} \neq 0, c_{1}=0\right\}$ and $\left\{\nu_{3}+v_{1}=0, \nu_{2} \neq 0, c_{2}=0\right\}$.

Conversely, if $\tilde{p}$ is a solution of (5.1) for some $\lambda<0, \mu$ and $\nu$, then we obtain a solution of (2.1) by setting $\tilde{p}_{0}=\tilde{p}(0), \tilde{p}_{T}=\tilde{p}(T), \tilde{r}_{0}=-\alpha \tilde{p}_{0}^{3}-\gamma$, and

$$
\begin{aligned}
& \tilde{v}_{1}=-\tilde{v}_{2}:= \pm \sqrt{2|\lambda|}, \\
& \tilde{c}_{1}=\frac{1}{2}\left(\frac{v}{\tilde{v}_{1}}-\frac{\mu}{2 \lambda}\right), \quad \tilde{c}_{2}=-\frac{1}{2}\left(\frac{v}{\tilde{v}_{1}}+\frac{\mu}{2 v}\right), \quad \tilde{c}_{3}=0 .
\end{aligned}
$$

The constant $\tilde{\nu}_{3}$ may be chosen arbitrarily.
A Painlevé II solution of the boundary value problem (2.1) may be used as a lower or an upper solution for a related boundary problem, for which Theorem 3.1 or Theorem 3.2 applies. In particular, the following result holds:

Corollary 5.1. Let $\tilde{p}$ be a nonnegative concave solution of (5.1) for some $\lambda<0, \mu \leqslant 0 \leqslant \nu$. Fix a constant $\tilde{v}_{3}>0$ and assume that $\tilde{p}(0)>0$. Then the boundary value problem (3.1) admits at least one solution $p$ such that $0 \leqslant p \leqslant \tilde{p}$ for any choice of the parameters for which:
(i) $0<p_{0} \leqslant \tilde{p}_{0}, 0 \leqslant p_{T} \leqslant \tilde{p}_{T}$.
(ii) $\nu_{1}+\nu_{2}+\nu_{3}=0, c_{1}+c_{2}+c_{3} \leqslant 0$.
(iii) $0 \leqslant \nu_{1} c_{1}+\nu_{2} c_{2}-c_{3}\left(\nu_{1}+\nu_{2}\right) \leqslant \nu$.
(iv) $-2 \lambda \leqslant \nu_{1}^{2}+v_{1} \nu_{2}+v_{2}^{2}$.
(v) $0<\nu_{1} \nu_{2}\left(\nu_{1}+\nu_{2}\right) \leqslant-2 \lambda \tilde{\nu}_{3}$.
(vi) $\tilde{\nu}_{3} \nu \leqslant \nu_{1} \nu_{2}\left(\nu_{1}+\nu_{2}\right)\left(c_{1}+c_{2}+c_{3}\right)$.
(vii) $\begin{aligned} \frac{\tilde{r}_{0}+v}{\tilde{p}_{0}}+\lambda \tilde{p}_{0}^{2}+\left(\tilde{v}_{3} v-\mu\right) j \leqslant & \frac{r_{0}+v_{1} c_{1}+v_{2} c_{2}-c_{3}\left(v_{1}+v_{2}\right)}{p_{0}}-\left(v_{1}^{2}+v_{1} \nu_{2}+v_{2}^{2}\right) \frac{p_{0}^{2}}{2} \\ & +\left[v_{1}^{2} c_{1}+v_{2}^{2} c_{2}-\left(\nu_{1}+v_{2}\right)^{2} c_{3}\right] j, \quad j=0,1 .\end{aligned}$

From conditions (ii) and (iii) it follows that $\tilde{C}_{3} \geqslant C_{3} \geqslant 0 . B y$ (iv), $\alpha \equiv 0$ is a lower solution of (3.1). Moreover, it follows from (vii) that $\tilde{\varphi} \leqslant \varphi$, and as $\tilde{C}_{1}=\tilde{v}_{3} \geqslant 0$ and $\tilde{p}$ is concave and nonnegative, then $\tilde{C}_{1}\left(\tilde{p}^{\prime}-\tilde{p}^{\prime}(0)\right) \tilde{p} \leqslant 0$. From (iv), we also deduce that $\tilde{C}_{2} \leqslant C_{2}$. Finally, from (ii), (v) and (vi) we conclude that $\tilde{C}_{4} \leqslant C_{4} \leqslant 0$ and $\tilde{C}_{5} \leqslant C_{5} \leqslant 0$. This implies that $\tilde{p}$ is an upper solution of (3.1).

It is noted that if $\tilde{v}_{3} \gg 0$, then conditions (ii) to (vi) are fulfilled for appropriate choices of $v_{i}$ and $c_{i}$. Moreover, condition (vii) holds if $r_{0} / p_{0}$ is large enough.

## References

[1] H.R. Leuchtag, J. Math. Phys. 22 (1981) 1317-1320.
[2] C. Rogers, A.P. Bassom, W.K. Schief, J. Math. Anal. Appl. 240 (1999) 367-381.
[3] C. De Coster, P. Habets, Two-Point Boundary Value Problems: Lower and Upper Solutions, Math. Sci. Eng., vol. 205, Elsevier, Amsterdam, 2006.
[4] P. Amster, M.C. Mariani, C. Rogers, C.C. Tisdell, J. Math. Anal. Appl. 289 (2004) 712-721.


[^0]:    * Corresponding author.

    E-mail address: pamster@dm.uba.ar (P. Amster).

